

Scaling limits of random fields with long range dependence¹

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Introduction

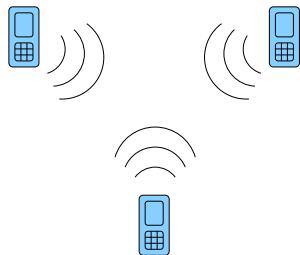
Outline

- ▶ Motivation
- ▶ Poisson noise model
- ▶ High-density limits
- ▶ Proofs

Earlier work:

- ▶ Taqqu and Levy (1986)
- ▶ Mikosch, Resnick, Rootzén, and Stegeman (2002)

Spatial noise generated by independent radio sources



- ▶ (X_j, R_j) location and transmission range of source j
- ▶ C shape of the transmission area
- ▶ λ mean density of the sources
- ▶ ρ mean transmission range

Describe the aggregate noise field

$$J_{\lambda, \rho}(x) = \# \{j : x \in X_j + R_j C\}$$

as $\lambda \rightarrow \infty$ and $\rho \rightarrow 0$.

Poisson grain field

Random scalar field on \mathbb{R}^d

$$J_{\lambda,\rho}(x) = \sum_{j=1}^{\infty} 1(x \in G_j),$$

generated by the *grains*

$$G_j = X_j + (\rho V_j)^{1/d} C$$

- ▶ deterministic shape: $C \subset \mathbb{R}^d$ bounded with unit volume
- ▶ random locations: $X_j \in \mathbb{R}^d$ independently scattered
- ▶ random volumes: $V_j \geq 0$ i.i.d. with unit mean
- ▶ model parameters: mean density λ and mean volume ρ

Poisson grain field in dimension one

For $C = [0, 1)$ in dimension one,

$$\begin{aligned} J_{\lambda, \rho}(t) &= \sum_{j=1}^{\infty} 1(t \in X_j + (\rho V_j)^{1/d} C) \\ &= \sum_{j=1}^{\infty} 1(X_j \leq t < X_j + \rho V_j) \end{aligned}$$

This is the $M/G/\infty$ queueing process:

- ▶ arrival instants X_j Poisson with intensity λ
- ▶ service times ρV_j i.i.d. with mean ρ

Heavy tails and regular variation

Assume

- ▶ $P(V_j > v) = L(v)v^{-\gamma}$
- ▶ $L(v)$ slowly varying:

$$\lim_{v \rightarrow \infty} \frac{L(\theta v)}{L(v)} = 1, \quad \text{for all } \theta > 0$$

Heavy tails when $\gamma \in (1, 2)$:

- ▶ $E V_j < \infty$
- ▶ $E V_j^2 = \infty$

Examples

- ▶ slowly varying: any $f(v)$ with $\lim_{v \rightarrow \infty} f(v) > 0$
- ▶ not slowly varying: $f(v) = e^{-v}$, $f(v) = 2 + \sin(v)$

High-intensity $M/G/\infty$ process with heavy tails:

(I) Long service times

Assume $P(V_j > v) = L(v)v^{-\gamma}$ with $\gamma \in (1, 2)$

Theorem (Mikosch, Resnick, Rootzén, Stegeman, 2002)

Assume $\lambda \rightarrow \infty$ and $\rho \rightarrow 0$ such that $\lambda\rho^\gamma L(1/\rho) \rightarrow \infty$. Then

$$\int_0^t \left(\frac{J_{\lambda,\rho}(s) - \mathbb{E} J_{\lambda,\rho}(s)}{c_{\lambda,\rho}} \right) ds \xrightarrow{fd} B_H(t).$$

B_H is fractional Brownian motion with $H = (3 - \gamma)/2$:

- ▶ Gaussian process
- ▶ $\mathbb{E} B_H(t) = 0$
- ▶ $\mathbb{E} B_H(s)B_H(t) = \frac{1}{2}(t^{2H} + s^{2H} - |t - s|^{2H})$

High-intensity $M/G/\infty$ process with heavy tails:

(II) Short service times

Theorem (Mikosch et al., 2002)

Assume $\lambda \rightarrow \infty$ and $\rho \rightarrow 0$ such that $\lambda \rho^\gamma L(1/\rho) \rightarrow 0$. Then

$$\int_0^t \left(\frac{J_{\lambda,\rho}(s) - \mathbb{E} J_{\lambda,\rho}(s)}{c_{\lambda,\rho}} \right) ds \xrightarrow{fd} \Lambda_\gamma(t).$$

Λ_γ is γ -stable Lévy process with unit skewness:

- ▶ independent increments
- ▶ γ -stable marginals with

$$\mathbb{E} e^{i\theta \Lambda_\gamma(t)} = e^{-t|\theta|^\gamma(1 - i \operatorname{sgn}(\theta) \tan(\pi\gamma/2))}$$

Poisson grain fields in higher dimensions

For dimensions $d > 1$,

$$\lim_{\lambda \rightarrow \infty, \rho \rightarrow 0} J_{\lambda, \rho}(x) = ?$$

Problems:

- ▶ The limit functions do not have point values (cf. derivative of Brownian motion)
- ▶ The cumulative process $\int_0^r J_{\lambda, \rho}(s\theta) ds$ does not characterize spatial correlation

Solution:

- ▶ Consider the *random linear functional*

$$J_{\lambda, \rho}(\phi) = \int_{\mathbb{R}^d} J_{\lambda, \rho}(x) \phi(x) dx$$

- ▶ For example for $d = 1$ with $\phi = 1_{[0, t]}$,

$$J_{\lambda, \rho}(\phi) = \int_0^t J_{\lambda, \rho}(s) ds$$

Poisson integral representation of $J_{\lambda,\rho}(x)$

$$\begin{aligned} J_{\lambda,\rho}(x) &= \sum_{j=1}^{\infty} 1(x \in X_j + (\rho V_j)^{1/d} C) \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}_+} 1(x \in y + v^{1/d} C) N_{\lambda,\rho}(dy, dv), \end{aligned}$$

- ▶ $N_{\lambda,\rho}$ is a Poisson random measure on with intensity

$$\mathbb{E} N_{\lambda,\rho}(\phi) = \lambda \int_{\mathbb{R}^d} \int_{\mathbb{R}_+} \phi(x, v) dx P(\rho V \in dv)$$

Poisson integral representation of $J_{\lambda,\rho}(\phi)$

- ▶ $J_{\lambda,\rho}$ as random linear functional

$$\begin{aligned} J_{\lambda,\rho}(\phi) &= \int_{\mathbb{R}^d} J_{\lambda,\rho}(x) \phi(x) dx \\ &= \int_{\mathbb{R}^d} \int_{\mathbb{R}_+} v m_\phi(x, v) N_{\lambda,\rho}(dx, dv) \end{aligned}$$

- ▶ m_ϕ are the averages

$$m_\phi(x, v) = \frac{\int_{v^{1/d}C} \phi(x+y) dy}{\int_{v^{1/d}C} dy}$$

High-density limit for light-tailed volumes

Theorem

Assume $E V^2 < \infty$. Then, as $\lambda \rightarrow \infty$ and $\rho \rightarrow 0$,

$$\frac{J_{\lambda,\rho}(\phi) - E J_{\lambda,\rho}(\phi)}{\rho(\lambda E V^2)^{1/2}} \xrightarrow{d} W(\phi), \quad \phi \in L^1 \cap L^2.$$

W is the *white Gaussian noise* on \mathbb{R}^d :

- ▶ Gaussian random linear functional
- ▶ $E W(\phi) = 0$
- ▶ $E W(\phi) W(\psi) = \int_{\mathbb{R}^d} \phi(x) \psi(x) dx$

Note that for $d = 1$:

- ▶ $W(1_{[0,t]}) = \int_0^t W(ds) = B(t)$
- ▶ $B(t)$ is Brownian motion

Grains with heavy-tailed volumes

- ▶ Assume $P(V > v) = v^{-\gamma}L(v)$ for $\gamma \in (1, 2)$
- ▶ Then for small ρ ,

$$\begin{aligned} E \sum_{j=1}^{\infty} 1(x \in G_j) 1(|G_j| > 1) &= \lambda \rho \int_{1/\rho}^{\infty} v P(V \in dv) \\ &\sim \frac{\lambda \rho^{\gamma} L(1/\rho)}{1 - \gamma^{-1}} \end{aligned}$$

- ▶ Three asymptotical regimes:

Small-grain scaling	$\lambda \rho^{\gamma} L(1/\rho) \rightarrow 0$
Intermediate scaling	$\lambda \rho^{\gamma} L(1/\rho) \rightarrow \gamma^{-1}$
Large-grain scaling	$\lambda \rho^{\gamma} L(1/\rho) \rightarrow \infty$

Three heavy-tailed limits: (I) Small-grain scaling

Theorem

If $\lambda \rightarrow \infty$ and $\rho \rightarrow 0$ such that $\lambda \rho^\gamma L(1/\rho) \rightarrow 0$, then

$$\frac{J_{\lambda,\rho}(\phi) - \mathbb{E} J_{\lambda,\rho}(\phi)}{\rho(1/(1-F))^\leftarrow(\gamma\lambda)} \xrightarrow{d} \Lambda_\gamma(\phi), \quad \phi \in L^1 \cap L^2.$$

Λ_γ is the independently scattered γ -stable random linear functional on \mathbb{R}^d characterized by

$$\mathbb{E} e^{i\Lambda_\gamma(\phi)} = e^{-\sigma_\phi^\gamma (1 - i\beta_\phi \tan(\frac{\pi\gamma}{2}))},$$

where $\sigma_\phi = \|\phi\|_\gamma$, $\beta_\phi = \|\phi\|_\gamma^{-\gamma} (\|\phi_+\|_\gamma^\gamma - \|\phi_-\|_\gamma^\gamma)$, and

$$\|\phi\|_\gamma = c_{\gamma,d} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \frac{\phi(x)\phi(y)}{|x-y|^{(1-\gamma)d}} dx dy$$

Three heavy-tailed limits: (II) Intermediate scaling

Theorem

If $\lambda \rightarrow \infty$ and $\rho \rightarrow 0$ such that $\lambda \rho^\gamma L(1/\rho) \rightarrow \gamma^{-1}$, then

$$J_{\lambda,\rho}(\phi) - \mathbb{E} J_{\lambda,\rho}(\phi) \xrightarrow{d} J_{\gamma,\mathcal{C}}^*(\phi), \quad \phi \in L^1 \cap L^2.$$

$J_{\gamma,\mathcal{C}}^*$ is the random linear functional on $L^1 \cap L^2$ with

$$J_{\gamma,\mathcal{C}}^*(\phi) = \int_{\mathbb{R}^d} \int_0^\infty v m_\phi(x, v) (N_\gamma(dx, dv) - dx v^{-\gamma-1} dv),$$

where $N_\gamma(dx, dv)$ is a PRM on $\mathbb{R}^d \times \mathbb{R}_+$ with intensity $dx v^{-\gamma-1} dv$

Three heavy-tailed limits: (III) Large-grain scaling

Theorem

If $\lambda \rightarrow \infty$ and $\rho \rightarrow 0$ such that $\lambda \rho^\gamma L(1/\rho) \rightarrow \infty$, then

$$\frac{J_{\lambda,\rho}(\phi) - \mathbb{E} J_{\lambda,\rho}(\phi)}{(\gamma \lambda \rho^\gamma L(1/\rho))^{1/2}} \xrightarrow{d} W_{\gamma,C}(\phi), \quad \phi \in L^1 \cap L^2.$$

$W_{\gamma,C}$ is the Gaussian random linear functional on $L^1 \cap L^2$

- ▶ $\mathbb{E} W_{\gamma,C}(\phi) = 0$
- ▶ $\mathbb{E} W_{\gamma,C}(\phi) W_{\gamma,C}(\psi) = \iint \phi(x) K_{\gamma,C}(x-y) \psi(y) dx dy$,
where

$$K_{\gamma,C}(x) = \int_0^\infty \left| v^{-1/d} x + C \cap C \right| v^{-\gamma} dv$$

Properties of the limits

Self-similarity

- ▶ A random linear functional J is *self-similar*, if

$$J(\phi \circ a^{-1}) \stackrel{d}{=} a^{Hd} J(\phi), \quad \text{for all } \phi \text{ and } a > 0$$

- ▶ The large-grain limit $W_{\gamma,C}$ is self-similar with $H = (3 - \gamma)/2$
- ▶ The small-grain limit Λ_γ is self-similar with $H = 1/\gamma$

Stationarity

- ▶ All limits $W_{\gamma,C}, J_{\gamma,C}^*, \Lambda_\gamma$ are stationary

Correlations

- ▶ $W_{\gamma,C}$ and $J_{\gamma,C}^*$ share the same second order statistics
- ▶ Λ_γ has infinite second moments

Fractional Gaussian noise

When the shape C is the closed unit ball,

$$W_{\gamma,C} \stackrel{d}{=} cW_H, \quad H = (3 - \gamma)/2,$$

where W_H is *fractional Gaussian noise* on $L^1 \cap L^2$:

▶ $E W_H(\phi) = 0$

▶ $E W_H(\phi) W_H(\psi) = \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \frac{\phi(x)\psi(y)}{|x-y|^{(2-2H)d}} dx dy$

The covariance kernel is called the *Riesz potential*

Fractional Gaussian noise in dimension $d = 1$

In dimension one,

$$W_H(1_{[0,t]}) = \int_0^t W_H(ds) = c_H B_H(t),$$

where B_H is fractional Brownian motion.

Hence

$$W_H = c_H \frac{d}{dt} B_H$$

in the sense of generalized functions.

White noise representation of W_H

For $H \in (1/2, 1)$,

$$W_H(\phi) \stackrel{d}{=} c_{H,d} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} \frac{\phi(y) dy}{|x - y|^{(3/2-H)d}} W(dx),$$

where W is white Gaussian noise on L^2

Proofs of Theorems 1–4

The proofs are based on

- (i) Fourier transform of the Poisson noise
- (ii) Extended Potter's bounds
- (iii) Hardy–Littlewood maximal function theory

Fourier transform of the Poisson noise

When N is a PRM with intensity η and

$$\int (|\phi| \wedge \phi^2) d\eta < \infty,$$

then the stochastic integral $\int \phi(dN - d\eta)$ exists, and

$$\mathbb{E} e^{i \int \phi(dN - d\eta)} = e^{\int \Psi(\phi) d\eta}$$

with

$$\Psi(v) = e^{iv} - 1 - iv.$$

Note: linearity \implies no need to consider $\sum_{j=1}^n \theta_j \phi_j$

Potter's bounds

Lemma

Let $V \geq 0$ be a random variable such that $P(V > v) = L(v)v^{-\gamma}$ for some $\gamma > 0$. Then for any $\epsilon \in (0, \gamma)$ there exist positive numbers c_ϵ and ρ_ϵ such that for all $\rho \in (0, \rho_\epsilon)$,

$$\frac{L(v/\rho)}{L(1/\rho)} \leq c_\epsilon (v^{-\epsilon} \vee v^{+\epsilon}) \quad \forall v > 0.$$

Hardy–Littlewood maximal theorem

Let ϕ_* be the *Hardy–Littlewood maximal function* of ϕ ,

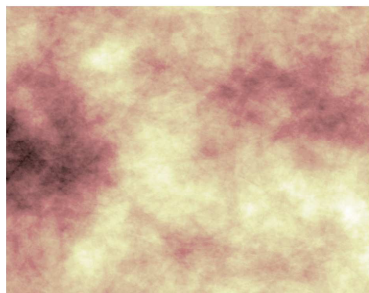
$$\phi_*(x) = \sup_{v>0} \frac{\int_{v^{1/d}C} |\phi(x+y)| dy}{\int_{v^{1/d}C} dy}$$

For all $p > 1$,

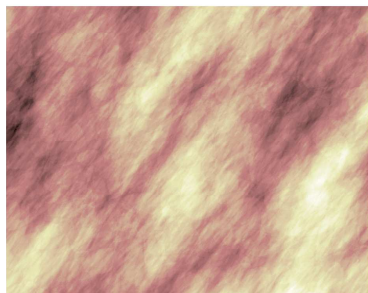
$$\phi \in L^p \implies \phi_* \in L^p$$

Illustration

$$J_{\lambda,\rho}(\phi) \sim E J_{\lambda,\rho}(\phi) + (\gamma\lambda\rho^\gamma)^{1/2} W_H(\phi)$$



Symmetric grains



Asymmetric grains

References

- ▶ M. Taqqu, J. B. Levy.
Using renewal processes to generate long-range dependence and high variability.
In *Dependence in Probability and Statistics*, pp. 73–89, Birkhäuser, 1986.
- ▶ T. Mikosch, S. Resnick, H. Rootzén, A. Stegeman.
Is network traffic approximated by stable Lévy motion or fractional Brownian motion?
Ann. Appl. Probab. **12**:23–68, 2002.
- ▶ I. Kaj, I. Norros, L. Leskelä, V. Schmidt.
Scaling limits of random fields with long-range dependence.
Ann. Probab. To appear, 2007.